

Supplementary Materials for “Superconducting micro-resonator arrays with ideal frequency spacing”

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Derivation of frequency collision probability and array yield formula

In this note, we assume an array of infinite number of resonators with the design resonance frequencies distributed in a geometric series, $f_n/f_{n-1} = 1 + \Delta$ ($\Delta \ll 1$, and $\Delta = 0.002$ for the array in our paper). We pick an arbitrary resonator in the array (referred to as the 0th-resonator) and assume there exist infinity number of resonators both above and below the frequency of the 0th resonator f_0 . This is an appropriate assumption for calculating the average yield for a large array of resonators. We model the actual resonance frequency of the n^{th} resonator as a random variable \tilde{f}_n with a mean of $\langle \tilde{f}_n \rangle = f_n$ and assume its fractional frequency error (between the design and actual resonance frequency), $\frac{\tilde{f}_n - f_n}{f_n}$, follows a Gaussian distribution of $G(0, \sigma^2/2)$ which is independent for different n .

We define a new random variable \tilde{x}_n by

$$\tilde{x}_n = \log\left(\frac{\tilde{f}_n}{f_0}\right) = \log(\tilde{f}_n) - \log(f_0), \quad (\text{S1})$$

which transforms the frequency axis into a log-frequency axis and fixes \tilde{x}_0 at the origin (no longer a random variable). The advantage of this transformation is that $x_n = \langle \tilde{x}_n \rangle$ is now evenly spaced by Δ on the x-axis ($x_n = n\Delta$, see Fig. S1), which simplifies our following analysis. It follows that \tilde{x}_n for each n also obeys an independent Gaussian distribution of $G(n\Delta, \sigma^2)$ (as shown in Fig. S1), where we have used the property that the difference of two independent Gaussian random variables is still a Gaussian random variable.

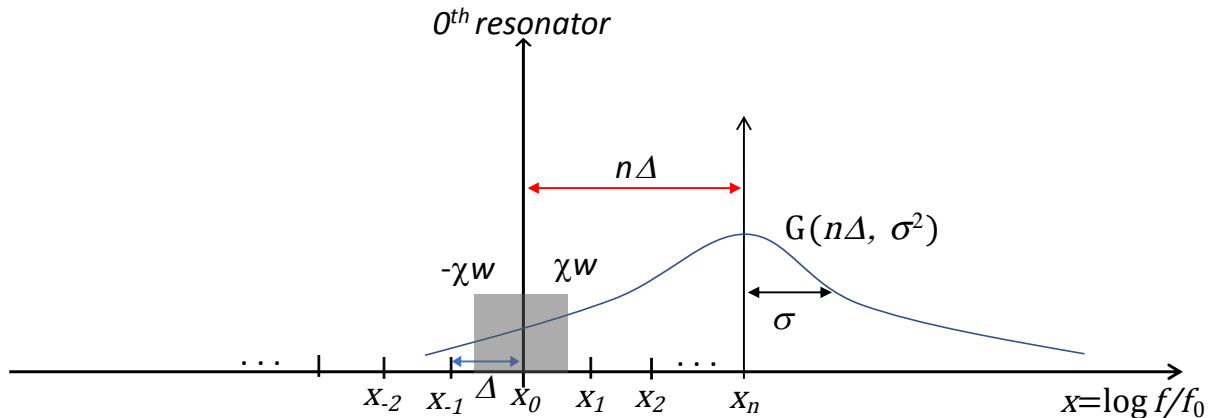


FIG. S1: Frequency collision model in log-frequency space.

We are now ready to calculate the probability P_{n0} for the 0th resonator to survive the collision with the n^{th} resonator ($n = \pm 1, \pm 2, \dots$). According to our χ -linewidth frequency

collision criterion, collision between the n^{th} and 0^{th} resonators occurs when \tilde{x}_n lies in the region of $[-\chi w, \chi w]$, where $w = 1/Q$ is the linewidth (unitless) of the resonator.

It follows that P_{n0} can be calculated in terms of an integral of the Gaussian distribution function in the region excluding $[-\chi w, \chi w]$,

$$\begin{aligned}
P_{n0} &= 1 - \int_{-\chi w}^{+\chi w} G(x; n\Delta, \sigma^2) dx & (S2) \\
&= 1 - \int_{\frac{-n\Delta - \chi w}{\sigma}}^{\frac{-n\Delta + \chi w}{\sigma}} G(x; 0, 1) dx \\
&= 1 - \frac{\text{Erf}\left(\frac{n\Delta + \chi w}{\sqrt{2}\sigma}\right) - \text{Erf}\left(\frac{n\Delta - \chi w}{\sqrt{2}\sigma}\right)}{2}
\end{aligned}$$

where Erf(x) is the error function defined as

$$\text{Erf}(x) = \frac{1}{\sqrt{\pi}} \int_{-x}^x e^{-t^2} dt \quad (S3)$$

After getting the probability for the 0_{th} resonator to survive the collision with the n_{th} resonator, we can calculate the probability for the 0_{th} resonator to survive collisions with all the other resonators ($n = \pm 1, \pm 2, \dots$) as

$$P_0 = \prod_n P_{n0} = \left\{ \prod_{n=1}^{n=\infty} \left[1 - \frac{\text{Erf}\left(\frac{n\Delta + \chi w}{\sqrt{2}\sigma}\right) - \text{Erf}\left(\frac{n\Delta - \chi w}{\sqrt{2}\sigma}\right)}{2} \right] \right\}^2 \quad (S4)$$

Eqn. S4 is the average yield of the array in terms of frequency collision.